Two-dimensional simulations of neutron star spreading layers

Pavel Abolmasov, Joonas Nättilä, Juri Poutanen

Tuorla Observatory, University of Turku

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Motivation for the study: importance of the BL

"Atoll" and "Z" sources



Variable contribution from the BL



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Motivation for the study: importance of the BL



Multiple QPOs at low (Hz to hHzs) and high (\sim 1kHz) frequencies



Sco X-1; van der Klis (any year, especially 2000)

nan

Motivation for the study: kHz QPOs and correlation time scales

kHzQPO frequency as a function of flux:



FIG. 1.—Left: Lower kHz QPO frequency vs. count rate in 4U 1608 – 52, showing the "parallel lines" phenomenon as it occurs in a single source, after Méndez et al. (1999). Right: kHz QPO frequencies vs. X-ray luminosity in 13 different sources, showing the parallel lines phenomenon across sources, after Ford et al. (2000).

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van der Klis (2006)

Motivation summary

- BL is visible
- BL is variable
- BL retains its structure (f_{kHzQPO} flux relation) on time scales minutes to hours, much longer than the scales of the inner disc or local thermal time scales
- ► ⇒spreading layer weakly coupled with both the disc and the NS

Spreading layer approximation

Spreading layer:



Inogamov& Sunyaev (1999, 2010), Suleimanov & Poutanen (2006) Need to set mixing/friction by hand or go to 2D/3D. But:

- surface friction should be small
- BC or matter source at the equator?

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BL simulations



Belyaev, Rafikov, and Stone (2012) ; Belyaev and Quataert (2018) Velocity discontinuity \Rightarrow shear instability. Inefficient in BL-disk angular momentum exchange. Rθ



2D spectral approach

Simulation setup:

- Decomposition into spherical harmonics (shtns library for python)
- pseudo-spectral approach with filtering
- https://github.com/ pabolmasov/SLayer
- goals: trace angular momentum transfer (shear instability!) and variability patterns

System of equations: continuity

$$\frac{\partial \Sigma}{\partial t} = -\nabla \cdot (\Sigma \boldsymbol{v}) + \boldsymbol{S}^{+} - \boldsymbol{S}^{-}, \qquad (1)$$

vorticity

$$\frac{\partial \omega}{\partial t} + \nabla \cdot (\omega \boldsymbol{v}) = -\nabla \times \frac{\nabla \Pi}{\Sigma} + \text{source/sink}, \quad (2)$$

divergence

$$\frac{\partial \delta}{\partial t} = \left[\nabla \times (\omega v)\right]_r - \nabla^2 \left(\frac{v^2}{2} + \Delta \Phi\right) - \nabla \cdot \left(\frac{1}{\Sigma} \nabla \Pi\right) + \text{sou}$$
(3)

energy

$$\frac{\partial E}{\partial t} + \nabla \cdot (Ev) = -\delta \Pi + Q^+ - Q^- + Q_{\rm NS} + Q_{\rm acc}.$$
(4)

Sub-sonic case

- KH stability depends on velocity gradients
- sonic modes
- heating where different velocities are mixed ⇒sometimes RT unstable
- rapid rotation near the SL boundary

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time-latitude diagram for angular frequency

Heating instability (setup-dependent)

nac

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super-sonic case, Keplerian rotation, $\dot{M} = 10^{-8} \text{ M}_{\odot} \text{ yr}^{-1}$



time-latitude diagram for surface density

Shear instability and oblique waves

Enhanced accretion solution ($\Sigma_0 = 10^8 \text{g cm}^{-2}$, $\dot{M} = 10^{-3} \text{ M}_{\odot} \text{ yr}^{-1}$) Reynolds's stress:



Power density spectra



1D-model

(net angular momentum conservation is still an option...) Angular momentum transfer:

$$\dot{M}j = 4\pi R^2 \sin^2 \theta R_{\theta\varphi} + \dot{M}j_0, \quad (5)$$

Energy balance:

$$\frac{cg_{\rm eff}}{\varkappa} (1 - \beta) = R_{\theta\varphi} \frac{\partial \Omega}{\partial \theta}.$$
 (6)
$$\Rightarrow \frac{\partial \omega}{\partial \theta} = \frac{(1 - \beta)\sin^2 \theta}{\dot{m}} \frac{1 - \omega^2 \sin^2 \omega}{\omega_0 - \omega \sin^2 \theta}.$$
 (7)



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Conclusions

- spectral simulations in a super-sonic regime are challenging but possible and provide a useful framework for further applications
- shear supersonic instabilities allow angular momentum transfer inside the SL (Reynolds's stress)
- non-linear evolution leads to quasiperiodics

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▶ why 10⁸ g cm⁻²?